

Astrophysics I: Stars and Stellar Evolution

AST 4001

Alexander Heger^{1,2,3}

¹School of Physics and Astronomy
University of Minnesota

²Nuclear & Particle Physics, Astrophysics & Cosmology Group, T-2
Los Alamos National Laboratory

³Department of Astronomy and Astrophysics
University of California at Santa Cruz

Stars and Stellar Evolution, Fall 2008

Overview

- 1 Recap
 - Regimes of Stellar Evolution
- 2 Stellar Evolution
 - Regimes of Stellar Evolution
- 3 Theory of the Main Sequence
 - Recap: Observational Behavior
 - Scaling Laws

Regimes of Stellar Evolution

- Recall

$$P_c = \sqrt[3]{4\pi B_n} GM^{2/3} \rho_c^{4/3}$$

- for ideal gas, $P_c = K_0 \rho_c T_c$ and we obtain

$$\rho_c = \frac{K_0^3}{4\pi B_n^3 G^3} \frac{T_c^3}{M^2}$$

\Rightarrow parallel lines with $\log \rho_c = 3 \log T - 2 \log M + \text{const.}$

- for non-rel. degenerate gas $P_c = K_1 \rho_c^{5/3}$ we obtain

$$\rho_c = 4\pi \left(\frac{B_{1.5G}}{K_1} \right)^3 M^2$$

\Rightarrow parallel lines at $\log \rho_c = 2 \log M + \text{const.}$

Quiz Solution

Find a relation for relativistic degenerate gas.

for rel. degenerate (electron) gas

$$P_c = K_2 \rho_c^{4/3}$$

in

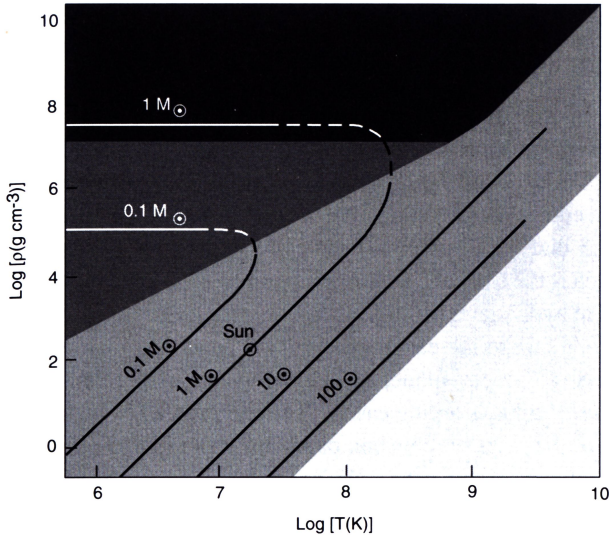
$$P_c = \sqrt[3]{4\pi} B_n G M^{2/3} \rho_c^{4/3}$$

we obtain (using $M_3 = (4B_3)^{-3/2}$)

$$M = \frac{1}{\sqrt{4\pi}} \left(\frac{K_2}{GB_3} \right)^{3/2} = 4\pi M_3 \left(\frac{K_2}{\pi G} \right)^{3/2}$$

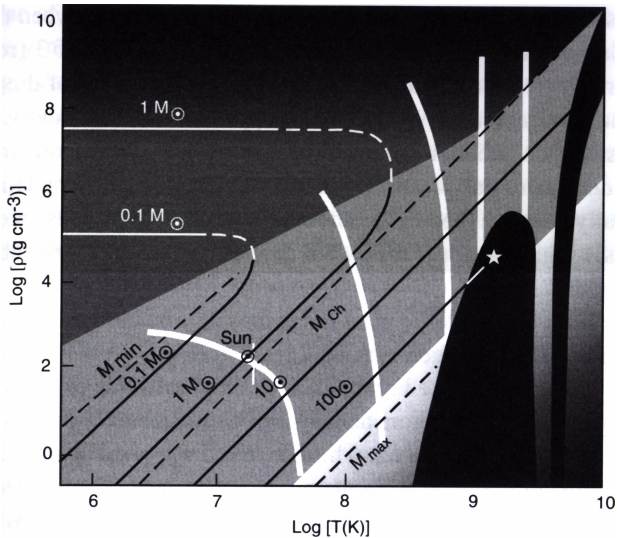
...the Chandrasekar Mass!

Domains of Stellar Mass



Regimes of stellar mass in the temperature-density diagram

Evolution Tracks



Evolution of
Stars in the
temperature-
density
diagram

Overview

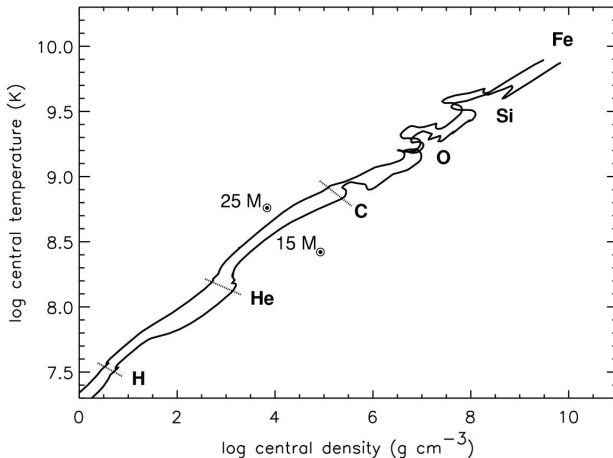
- 1 Recap
 - Regimes of Stellar Evolution
- 2 Stellar Evolution
 - Regimes of Stellar Evolution
- 3 Theory of the Main Sequence
 - Recap: Observational Behavior
 - Scaling Laws

Burning Phases in Stars

20 M_⊙ star

Fuel	Main Product	Secondary Product	T (10 ⁹ K)	Time (yr)	Main Reaction
H	He	¹⁴ N	0.02	10 ⁷	^{CNO} 4 H → ⁴ He
He	O, C	¹⁸ O, ²² Ne s-process	0.2	10 ⁶	3 He ⁴ → ¹² C ¹² C(α,γ) ¹⁶ O
C	Ne, Mg	Na	0.8	10 ³	¹² C + ¹² C
Ne	O, Mg	Al, P	1.5	3	²⁰ Ne(γ,α) ¹⁶ O ²⁰ Ne(α,γ) ²⁴ Mg
O	Si, S	Cl, Ar, K, Ca	2.0	0.8	¹⁶ O + ¹⁶ O
Si, S	Fe	Ti, V, Cr, Mn, Co, Ni	3.5	0.02	²⁸ Si(γ,α)...

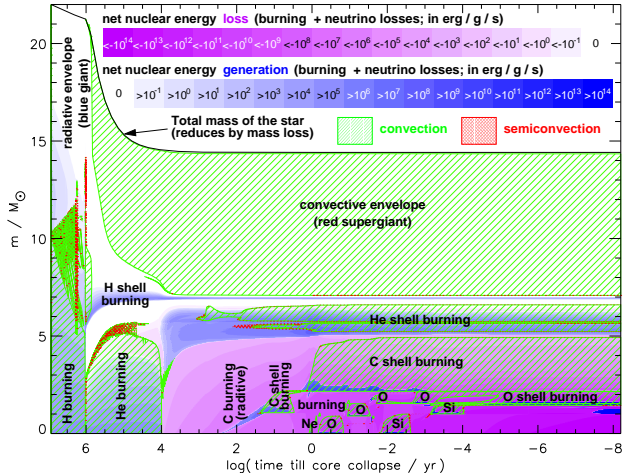
Evolution of Stars, $15 M_{\odot}$ and $25 M_{\odot}$



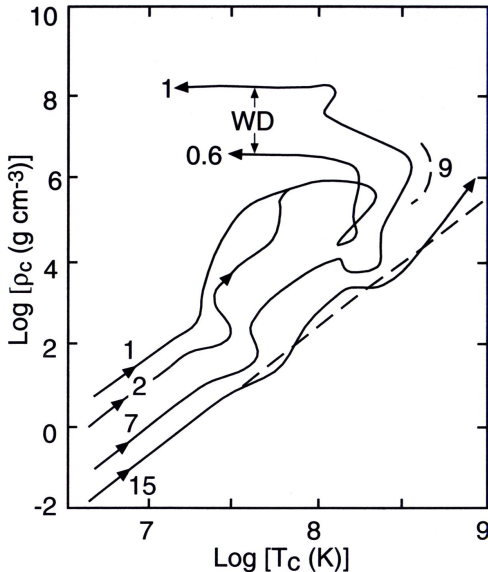
Evolution of central temperature and density for initial stellar masses of $15 M_{\odot}$ and $25 M_{\odot}$ in the density-temperature diagram

(note reversal of T and ρ)

Burning Phases in the Stellar Interior

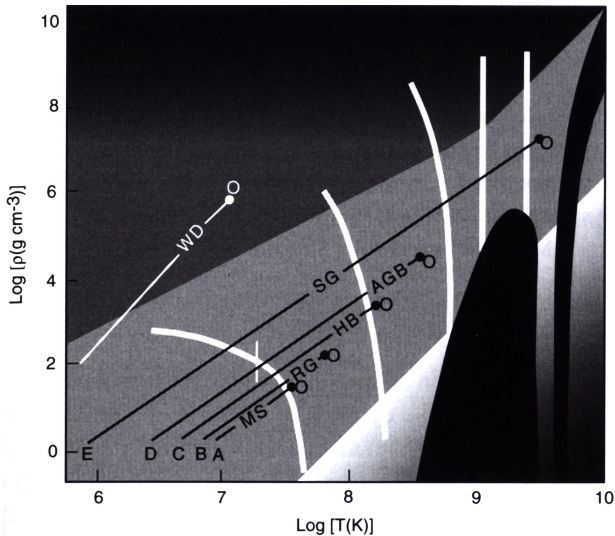


Evolution of Stars, 1–15 M_{\odot}



Evolution of central temperature and density for initial stellar masses of $1 M_{\odot}$ to $15 M_{\odot}$ in the temperature-density diagram.

Configuration of a $10 M_{\odot}$ Star

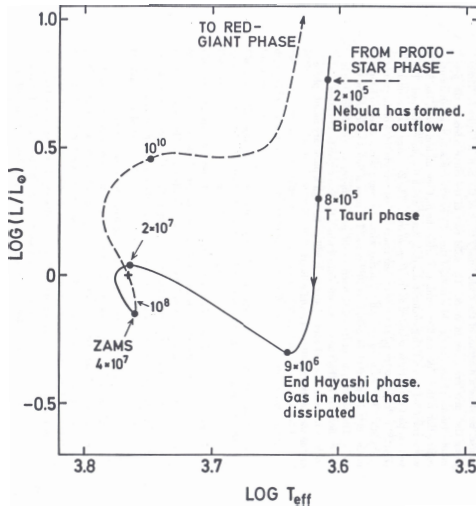


Configuration of a $10 M_{\odot}$ star at different evolution phases in the temperature-density diagram.

Overview

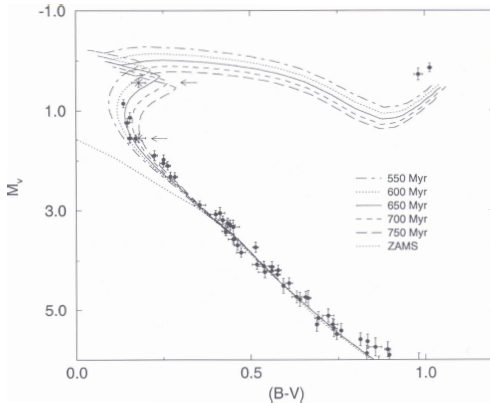
- 1 Recap
 - Regimes of Stellar Evolution
- 2 Stellar Evolution
 - Regimes of Stellar Evolution
- 3 Theory of the Main Sequence
 - Recap: Observational Behavior
 - Scaling Laws

Evolution of the Sun in the HRD



Evolution of the sun
from formation
through hydrogen
burning

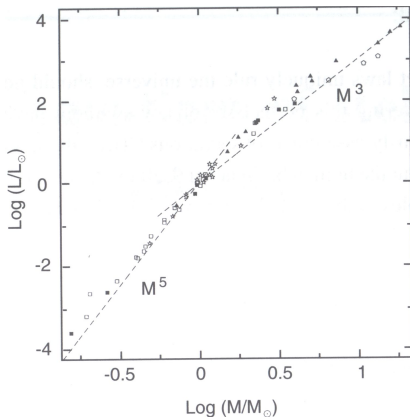
Main Sequence in a Star Cluster



Hyades cluster and stellar tracks

$$\log L = \alpha \log T_{\text{eff}} + \text{const.}$$

Mass-Luminosity Relation



Mass-Luminosity relation for
(zero-age) main-sequence
(ZAMS) stars

$$L \propto M^\nu$$

with $\nu = 3 \dots 5$.

Can be calibrated piecewise to

$$\frac{L}{L_\odot} = \left(\frac{M}{M_\odot} \right)^\nu$$

Stellar Structure Equations

stationary terms

time-dependent terms

$$\frac{\partial r}{\partial m} = \frac{1}{4\pi r^2 \rho} \quad (1)$$

$$\frac{\partial P}{\partial m} = -\frac{Gm}{4\pi r^4} - \frac{1}{4\pi r^2} \frac{\partial^2 r}{\partial t^2} \quad (2)$$

$$\frac{\partial F}{\partial m} = \varepsilon_{\text{nuc}} - \varepsilon_{\nu} - c_P \frac{\partial T}{\partial t} + \frac{\delta}{\rho} \frac{\partial \rho}{\partial t} \quad (3)$$

$$\frac{\partial T}{\partial m} = -\frac{GmT}{4\pi r^4 P} \nabla \left[1 + \frac{r^2}{Gm} \frac{\partial^2 r}{\partial t^2} \right] \quad (4)$$

$$\frac{\partial X_i}{\partial t} = f_i(\rho, T, \mathbf{X}) \quad (5)$$

where $\mathbf{X} = \{X_1, X_2, \dots, X_i, \dots\}$.

Simplified Stellar Structure Equations

- only radiative temperature gradient: $\nabla = \nabla_{\text{rad}} = \frac{3}{16\pi acG} \frac{\kappa FP}{mT^4}$
- only simple law for nuclear burning: $\varepsilon = \varepsilon_{\text{nuc}} - \varepsilon_{\nu} = q_0 \rho T^n$
- only ideal gas pressure: $P = P_{\text{gas}} = \frac{\mathcal{R}T\rho}{\mu}$

$$\begin{aligned}\frac{\partial r}{\partial m} &= \frac{1}{4\pi r^2 \rho} \\ \frac{\partial P}{\partial m} &= -\frac{Gm}{4\pi r^4} \\ \frac{\partial T}{\partial m} &= -\frac{3}{4ac} \frac{\kappa}{T^3} \frac{F}{(4\pi r^2)^2} \\ \frac{\partial F}{\partial m} &= q_0 \rho T^n \\ P &= \frac{\mathcal{R}T\rho}{\mu}\end{aligned}$$

Dimensionless Stellar Structure Equations

- all functions $r(m)$, $P(m)$, $\rho(m)$, $T(m)$, and $F(m)$ need to be solved in range $0 \leq m \leq M$
- free parameter: mass M
- parameters κ , q_0 , μ , and n determined from physics
- introduce dimension-less variable

$$x = \frac{m}{M}$$

- we can now write a set of dimension-less equations with functions $f_i(x)$ for these 5 quantities:

$$r = f_1(x)R_*$$

$$P = f_2(x)P_*$$

$$\rho = f_3(x)\rho_*$$

$$T = f_4(x)T_*$$

$$F = f_5(x)F_*$$

Dimensionless Stellar Structure Equations

- substituting

$$m = Mx,$$

$$r = f_1(x)R_*,$$

$$P = f_2(x)P_* \text{ into}$$

$$\frac{\partial P}{\partial m} = -\frac{Gm}{4\pi r^4}$$

we obtain

$$\frac{P_*}{M} \frac{df_2}{dx} = -\frac{GMx}{4\pi f_1^4 R_*^4}$$

- If we define

$$P_* = \frac{GM^2}{R_*^4}$$

we may write

$$\frac{df_2}{dx} = -\frac{x}{4\pi f_1^4}$$

Quiz

Try this for the other four equations

- Work and discuss in groups of 2-3.
- 3 min
- Please write up your solution.
- Please sign with your names and to hand in.
- (no grades)

Quiz - equations

Try this for the other four equations

$$\frac{\partial r}{\partial m} = \frac{1}{4\pi r^2 \rho}$$

$$\frac{\partial P}{\partial m} = -\frac{Gm}{4\pi r^4}$$

$$\frac{\partial T}{\partial m} = -\frac{3}{4ac} \frac{\kappa}{T^3} \frac{F}{(4\pi r^2)^2}$$

$$\frac{\partial F}{\partial m} = q_0 \rho T^n$$

$$P = \frac{\mathcal{R} T \rho}{\mu}$$

$$r = f_1(x) R_*$$

$$P = f_2(x) P_*$$

$$\rho = f_3(x) \rho_*$$

$$T = f_4(x) T_*$$

$$F = f_5(x) F_*$$

Dimensionless Stellar Structure Equations

In a similar way we can re-write the entire set

$$\begin{aligned}\frac{df_2}{dx} &= -\frac{x}{4\pi f_1^4} & , & & P_* &= \frac{GM^2}{R_*^4} \\ \frac{df_1}{dx} &= \frac{1}{4\pi f_1^2 f_3} & , & & \rho_* &= \frac{M}{R_*^3} \\ f_2 &= f_3 f_4 & , & & T_* &= \frac{\mu P_*}{\mathcal{R} \rho_*} \\ \frac{df_4}{dx} &= -\frac{3f_5}{4f_4^3 (4\pi f_1^2)^2} & , & & F_* &= \frac{ac}{\kappa} \frac{T_*^4 R_*^4}{M} \\ \frac{df_5}{dx} &= f_3 f_4^n & , & & F_* &= q_0 \rho_* T_*^n M\end{aligned}$$

\Rightarrow *homology* of solution as function of M !

Dimensionless Stellar Structure Equations

- substituting $P_* = \frac{GM^2}{R_*^4}$ and $\rho_* = \frac{M}{R_*^3}$ into $T_* = \frac{\mu P_*}{\mathcal{R} \rho_*}$ we obtain:

$$T_* = \frac{\mu G M}{\mathcal{R} R_*}$$

- adding this into $F_* = \frac{ac}{\kappa} \frac{T_*^4 R_*^4}{M}$ we obtain

$$F_* = \frac{ac}{\kappa} \left(\frac{\mu G}{\mathcal{R}} \right)^4 M^3$$

- we recover $L \propto M^3$
- $\tau_{\text{MS}} = \frac{M}{L} \propto M^{-2}$
- but this relation also holds for any value of x -
inside star at same relative mass coordinate

Dimensionless Stellar Structure Equations

- substituting $F_* = \frac{ac}{\kappa} \left(\frac{\mu G}{\mathcal{R}} \right)^4 M^3$ into $F_* = q_0 \rho_* T_*^n M$
and using $P_* = \frac{GM^2}{R_*^4}$, $\rho_* = \frac{M}{R_*^3}$, and $T_* = \frac{\mu P_*}{\mathcal{R} \rho_*}$ we obtain:

$$R_* \propto M^{\frac{n-1}{n+3}}$$

- \Rightarrow for large n (CNO cycle: $n \approx 14 \dots 16$):
roughly $R_* \propto M$
(big stars)
- \Rightarrow for small n (pp chains: $n = 4$):
 $R \propto M^{3/7}$

Dimensionless Stellar Structure Equations

- inserting $R_* \propto M^{\frac{n-1}{n+3}}$ into $\rho_* = \frac{M}{R_*^3}$ we obtain

$$\rho_* \propto M^2 \frac{3-n}{3+n}$$

- since $n > 3$:
density decreases with mass!
- in particular, this is true for central density.

Dimensionless Stellar Structure Equations

- Using

$$L = 4\pi R^2 \sigma T_{\text{eff}}^4$$

we obtain

$$L^{1 - \frac{2(n-1)}{3(n+3)}} \propto T_{\text{eff}}^4$$

- for $n = 4$ we obtain

$$\log L = 5.6 \log T_{\text{eff}} + \text{const.}$$

- for $n = 16$ we obtain

$$\log L = 8.4 \log T_{\text{eff}} + \text{const.}$$

Dimensionless Stellar Structure Equations

- from $T_* = \frac{\mu G}{\mathcal{R}} \frac{M}{R_*}$ and $R_* \propto M^{\frac{n-1}{n+3}}$ we obtain for the (central) temperature

$$T_c \propto M^{\frac{4}{n+3}}$$

- for $n = 4$ (pp chain, low-mass stars) we hence have

$$T_c \propto M^{4/7}$$

- for $n = 16$ (CNO cycle, massive stars) we hence have

$$T_c \propto M^{1/5}$$

- \Rightarrow due to high temperature-sensitivity of CNO cycle nuclear burning, massive stars require only little higher central temperature to compensate for their high luminosity.

Dimensionless Stellar Structure Equations

- calibration to the sun, $T_{c,\odot} \approx 1.5 \times 10^7$ K

$$\frac{T_c}{T_{c,\odot}} = \left(\frac{M}{M_\odot} \right)^{4/7}$$

- assuming minimum temperature T_{\min} for hydrogen ignition, $T_{\min} \approx 4 \times 10^6$ K, and requiring $T_c > T_{\min}$ we obtain

$$\frac{M}{M_\odot} \geq \left(\frac{T_{\min}}{T_{c,\odot}} \right)^{7/4}$$

- \Rightarrow minimum stellar mass $M_{\min} \approx 0.1 M_\odot$
- \Rightarrow minimum stellar luminosity:

$$\frac{L_{\min}}{L_\odot} = \left(\frac{M_{\min}}{M_\odot} \right)^3 \approx 10^{-3}$$