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Nuclear Physics I: Nuclear Astrophysics PHYS 8801

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Nuclear Physics I: Nuclear Astrophysics, Spring 2012

Supernova	е

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Agenda

Supernovae

- Supernova Types and Light Curves
- 2 Black Holes
 - Kerr Black Holes

3 Binary Stars

- Binary Types
- The Roche Model
- Interacting Binaries



Supernova Types as Function of Mass and Metallicity

(single stars)

	SN Type	pre-SN stellar structure	
	llp	> 2 M $_{\odot}$ H envelo	pe
	IIL	< 2 M_{\odot} H envelo	ре
	lb/c	no H envelo	pe
Type lb/c He core mass at explosion		explosion energy	display
> 15 N	∕I₀	direct collapse	none
~158	M₀	weak	dim
~8…5 M _☉		strong	dim
< 5 N	l _o	strong	bright

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Supernovae o●oooo	Black Holes	Binary Stars	Nuclear Masses
Supernovae			

Sequence of increasingly stripped cc SNe



Supernovae		
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Black Holes

Binary Stars

Core Collapse Supernovae



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Core Collapse Supernovae – 3D

Cold inflow and **hot outflow** in 3D simulations → similar to dipolar flow pattern observed in 2D rotationally symmetric simulations



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Supernovae ooooooo	Black Holes	Binary Stars	Nuclear Ma
Supernovae			

Neutron Star Kicks



Dipolar oscillation may explain observed neutron star kicks of several 100 km/s.

Explosive Nucleosynthesis

in supernovae

Fuel	Main Product	Secondary Product	Т (10 ⁹ К)	Time (s)	Main Reaction
Innermost ejecta	<i>r</i> -process	-	>10 low Y _e	1	(n, γ), β ⁻
Si, O	⁵⁶ Ni	iron group	>4	0.1	(α,γ)
Ο	Si, S	CI, Ar, K, Ca	3 - 4	1	¹⁶ O + ¹⁶ O
O, Ne	O, Mg, Ne	Na, Al, P	2 - 3	5	(γ,α)
		p-process ¹¹ B, ¹⁹ F, ¹³⁸ La, ¹⁸⁰ Ta	2 - 3	5	(γ ,n)
		v-process		5	(v, v'), (v, e ⁻)

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Black Holes

Binary Stars

Nuclear Masses

Black Holes



Saturn as seen through the gravitational lense of a black hole

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Black Holes

Binary Stars

Nuclear Masses

Density and Radii of Astronomical Objects



- comparison of average density of astronomical objects
- average density of black holes decreases with increasing mass

 $R_{
m s}\sim M$ $V\sim R^{3}$ $ar{
ho}=M/V\sim M^{-2}$

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Black Holes

Binary Stars

Black Holes - Angular Momentum and Orbits



in this figure the following conventions were used:

- "gravitational radius"
 r_g = R_s
- normalized angular momentum $\tilde{L} = I/cmr_g$ where *I* is the angular momentum of the particle
- specific *total* energy of the particle $\tilde{E} = E/mc^2$ $E = mc^2 + E_{bind} + E_{kin}$

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Black Holes

Binary Stars

Black Holes - Energy and Orbits



Fig. 3. Effective black hole potential. $1 - \widetilde{E} = \widetilde{E}_1$, $2 - \widetilde{E} = \widetilde{E}_2$, $3 - \widetilde{E} = \widetilde{E}_3$, $4 - \widetilde{E} = \widetilde{E}_4$



Types of orbits

- 1/a: bound/closed, $0 < \tilde{E} < \tilde{E}_{max}$
- 2/b: unbound/open, $\tilde{E} < 0$
- 3/c: "unbound"/capture for $\tilde{E} > \tilde{E}_{max} > 0$ (not in classical mechanics)
- 4/d: closed capture loop $\tilde{E} < \tilde{E}_{max}$, can be $\tilde{E} < 0$ or $\tilde{E} > 0$ (not in classical mechanics)

Black Holes

Binary Stars

Nuclear Masses

Kerr Black Holes



Fig. 7. A rotating black hole: 1-horizon, 2-ergosphere, 3-static limit

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Black Holes

Binary Stars

Nuclear Masses

Particle Orbit around a Kerr Black Hole



Black Holes

Binary Stars

Nuclear Masses

Critical Radii in Black Holes

Orbit	a = 0	a = M	
		L > 0	L < 0
$r_{\rm photon}$	1.5	0.5	2.0
r _{bind}	2.0	0.5	2.92
$r_{\rm bound}$	3.0	0.5	4.5

Schwarzschild Case

photon circular orbit

$$\textit{r}_{bind} = 1.5\textit{R}_{s} = 3\frac{\textit{GM}}{\textit{c}^{2}}$$

last stable orbit

$$r_{\text{bound}} = 3R_{\text{s}} = 6\frac{GM}{c^2}, \ v = \frac{c}{2}$$

 last marginally stable circular orbit

$$\mathit{r}_{\mathsf{bind}} = 2\mathit{R}_{\mathsf{s}} = 4 \frac{\mathit{GM}}{\mathit{c}^2} \,, \, \mathit{v} = \mathit{v}_{\mathsf{esc}}$$

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Black Holes 0000000

Orbits and Energies



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Summary - Comparison of Compact Remnants

Distinguishing Traits of Compact Objects				
Object	Mass ^a (M)	Radius ^b (R)	Mean Density (g cm ⁻³)	Surface Potential (GM/Rc^2)
Sun White dwarf Neutron star Black hole	$M_{\odot} \leq M_{\odot} \\ \sim 1-3M_{\odot} \\ \text{Arbitrary}$	R_{\odot} $\sim 10^{-2}R_{\odot}$ $\sim 10^{-5}R_{\odot}$ $2GM/c^{2}$	$1 \\ \leq 10^7 \\ \leq 10^{15} \\ \sim M/R^3$	$ \begin{array}{r} 10^{-6} \\ \sim 10^{-4} \\ \sim 10^{-1} \\ \sim 1 \end{array} $

 ${}^{a}M_{\odot} = 1.989 \times 10^{33} \text{ g}$ ${}^{b}R_{\odot} = 6.9599 \times 10^{10} \text{ cm}$

Black Holes

Binary Stars

Nuclear Masses

Binary Stars ('Binaries')



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Black Holes

Binary Stars

Nuclear Masses

Roche Model

$$\begin{aligned}
\phi(x,y,z) &= -\frac{GM_{1}}{1r_{1}^{2}1} - \frac{GM_{2}}{1r_{2}^{2}1} - \frac{1}{2} |\vec{S}|^{2} \omega^{2} & \xrightarrow{M_{1}} y^{2} + \frac{y^{2}}{r_{2}} |\vec{r}_{2}| \\
& \text{centrifugal potential} \\
& |\vec{r}_{1}| &= (x^{2} + y^{2} + z^{2})^{1/2} , \quad |\vec{r}_{2}| &= ((A - x)^{2} + y^{2} + z^{2})^{1/2} \\
& |\vec{S}| &= ((x - x_{s})^{2} + y^{2})^{1/2} &= \left[\left(x - \frac{M_{2}}{M_{1} + M_{2}} A \right)^{2} + y^{2} \right]^{1/2} \\
& \omega^{2} &= \frac{G(M_{1} + M_{2})}{A^{3}} &: \quad 3^{rd} \text{ Kepler's Law} \\
& \text{Introduce dimensionless variables:} \quad \vec{\xi} &= \frac{x}{A} ; \quad \vec{\gamma} &= \frac{y}{A} ; \quad \vec{\xi} &= \frac{x}{A} ; \quad \vec{\gamma} &= \frac{y}{A} ; \quad \vec{\xi} &= \frac{x}{A} ; \quad \vec{\gamma} &= \frac{y}{A} ; \quad \vec{\xi} &= \frac{x}{A} ; \quad \vec{\gamma} &= \frac{y}{A} ; \quad \vec{\xi} &= \frac{x}{A} ; \quad \vec{\xi} &= \frac{x}{A$$

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Black Holes

Binary Stars

Nuclear Masses

Roche Potential



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Black Holes

Binary Stars

Nuclear Masses

Lagrange Points



Five Lagrange points:

L1, L2, L3: unstable

L4,L5: stable

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Black Holes

Binary Stars

Nuclear Masses

Contact Binaries



 Supernovae
 Black Holes
 Binary Stars
 Nuclear Masses

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Stability of mass transfer depends on reaction of donor and receiving star



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Star + compact remnant + Roche-Lobe overflow: X-ray binaries

WD + companion: Novae, Dwarf Novae, Type la supernovae

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NS + companion: X-ray bursts, millisecond pulsars, ...

NS+NS: Binary pulsars

Black Holes

Binary Stars

Binary Pulsar Production



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