

# NUCLEAR MASSES

## BINDING ENERGY

$$B(N, Z) = (N \cdot m_n + Z \cdot m_p - m(N, Z)) c^2$$

→ less than 1% of mass

$$B = -E$$

$$H\psi = E\psi$$

## ATOMIC MASSES

$$m_{\text{at}}(N, Z) = m(N, Z) + Z \cdot m_e - B_{\text{el}}(Z)/c^2$$

with

$$B_{\text{el}}(Z) \approx (14.4381 Z^{2.39} + 1.53468 \times 10^{-6} Z^{5.35}) \text{ eV}$$

$$\text{MASS EXCESS: } \Delta M(N, Z) = (m_{\text{at}}(N, Z) - \frac{A}{12} m_{\text{at}}(6, 6)) \times c^2$$

# SEMI-EMPIRICAL MASS FORMULA

WEIZSÄCKER (1935), BETHE & BACHER (1936)

$$E = a_{\text{vol}} \cdot A + a_{\text{SF}} \cdot A^{2/3} + \frac{3e^2}{5r_0} Z^2 A^{-1/3} \\ + (a_{\text{sym}} \cdot A + a_{\text{SS}} \cdot A^{2/3}) I^2 \\ + \Delta n_i + \Delta p$$

$\Delta n_i, p = \pm \delta$  for  $N, Z$  even or odd

(RMS charge Radius:  $R_{\text{ch}} = r_0 \cdot A^{1/3}$ )

$$I = \frac{N-Z}{A}$$

$a_{\text{SS}}$ : correction by Myers & Swiatec (1966)

$$a_{\text{vol}} = -15.697550 \text{ MeV}$$

$$a_{\text{SF}} = 17.662690 \text{ MeV}$$

$$a_{\text{sym}} = 26.308165 \text{ MeV}$$

$$a_{\text{SS}} = -17.003132 \text{ MeV}$$

$$r_0 = 1.221897 \text{ fm}$$

$$\delta = -1.250000 \text{ MeV}$$

$$e^2 = 1.43985 \text{ MeV} \cdot \text{fm}$$

## LARGE - A LIMIT OF WEIZSÄCKER FORMULA

$$E = \frac{E}{A} = a_{\text{vol}} + a_{\text{SF}} A^{-1/3} + \frac{3e^2}{20r_0} A^{2/3} (1-I) + (a_{\text{sym}} + a_{\text{as}} A^{-1/3}) I^2$$

$$\text{AGAIN } I = \frac{N-Z}{A}$$

COULOMB TERM: (NON-SATURATING)

→ diverge unless  $I = 1$

→ PURE neutron system

IF we set  $I = 1$ :

$$E \sim a_{\text{vol}} + a_{\text{sym}}$$

→  $\sim +10.6 \text{ MeV}$

Q: ARE LARGE  $n$ -systems unbound?

WHAT ABOUT NS?

→ BOUND BY GRAVITY!

→ DERIVE WEISSÄCKER FORMULA w/ GRAVITY!

(RECALL: IDENTICAL TO COULOMB  
BUT: ATTRACTIVE)

$$E_G = \frac{3GM^2}{5R} = \frac{3Gm_n^2}{5r_0} A^{2/3}$$

For const  $\rho$

$$e = a_{vol} + a_{sf} A^{-1/3} + \frac{3}{5r_0} \left\{ \frac{e^2}{4} (1-I)^2 - Gm_n^2 \right\} A$$

$$- (a_{sym} + a_{ss} A^{-1/3}) I^2$$

AGAIN: large system, COULOMB RE PUSION  
SET  $I = 1$

$$e \sim a_{vol} + a_{sym} - \frac{3Gm_n}{5v_0} A^{2/3}$$

for large  $A$   $Gm_n$  will eventually dominate

→ system will be bound

$$e < 0 \text{ for } A > \left\{ \frac{5v_0}{3Gm_n^2} (a_{vol} + a_{sym}) \right\}^{3/2}$$

$$\approx 0.179 \times 10^{56}$$

$$\text{typical } A_{NS} \approx 1.6 \times 10^{57}$$

→ NS within a factor 20

Realistic  $M_{DFC}$

$$M_{MIN, NS} \sim 0.1 M_{\odot} \sim 10^{56} \text{ neutrons}$$

But: 1) NS not pure n-matter:

$\beta$ -decays to p-e pairs

2) central density  $\rho_c \sim 4 \rho_0$

$\rightarrow \mu$  (neutrons) and maybe free quarks  
at center?

3) treat in GR

(consequence: BINARY NS:

ONE ACCRETS, secondary eventually  
falls below lower limit  
 $\rightarrow$  blow up!

$\rightarrow$  Symmetry energy blows up stars  
when grav. binding too low

$$E = E_s(N+Z) + E_n(N-Z)$$

Shell model

↔ single particle levels

$$\vec{j} = \vec{l} + \vec{s}$$

→ occupancy #  $2j+1$

Substates  $m_j = -j, \dots, j$

Parity  $\pi = (-1)^l$

Spectroscopic notation  $s, p, d, f, g, \dots$

for  $l = 0, 1, 2, 3, 4, \dots$

Spin-orbit tensor

$$j = l + \frac{1}{2} \text{ or } j = l - \frac{1}{2}$$

$l$  integer →  $j$   $\frac{1}{2}$ -integer

here: very strong spin-orbit coupling

→ significant change of single-particle states!

RESULT:

FORMATION OF ENERGY GAPS → SHELLS

→ magic #s: 2, 8, 20, 28, 50, 82, 128